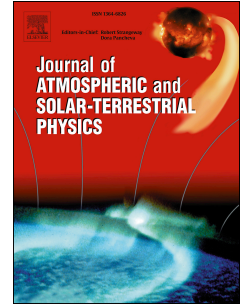


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PPEF Responses

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1 **IMF By-Driven Electric Field Disturbances Over the Equator During Northward**
2 **IMFs: PPEF Responses**

3
4 Ram Singh^{1,2}, Tarun Kumar Pant², S. Sripathi³, D. Rout⁴, A. Bhaskar², and Danny E. Scipion¹

5
6 ¹Radio Observatorio de Jicamarca, Instituto Geofísico del Perú, Lima, Peru

7 ²Space Physics Laboratory, Vikram Sarabhai Space Centre, Trivandrum, India

8 ³Indian Institute of Geomagnetism (IIG), New Panvel, Navi Mumbai, India

9 ⁴National Atmospheric Research Laboratory, Gadanki, India,

10
11 Corresponding author: Ram Singh (ramphysics4@gmail.com)

12
13
14 **Abstract**

15 This article provides the first evidence of equatorial and low-latitude ionospheric electric field
16 disturbances driven by quasi-periodic variations in the interplanetary magnetic field (IMF) By
17 component, under a constant northward IMF Bz ($\sim +15$ nT), solar wind dynamic pressure (~ 8
18 nPa), solar wind velocity (~ 450 km/s), and positive Sym-H (+10 nT). The virtual height of the
19 F layer ($h'F$), measured by ionosondes at opposite longitudes in the American (Jicamarca,
20 11.9°S , -76.0°E ; Fortaleza, 3.9°S , -38.52°E) and Indian (Trivandrum, 8.5°N , 77.0°E ;
21 Tirunelveli, 8.73°N , 77.7°E) sectors, shows electric field disturbances with opposite polarities
22 during the day and night, driven by changes in the IMF. During northward IMF Bz conditions,
23 vertical $E \times B$ plasma drifts measured by the Jicamarca Incoherent Scatter Radar (ISR), along
24 with equatorial electrojet (EEJ) observations, reveal westward electric field disturbances on the
25 dayside. Conversely, on the nightside, ionosonde measurements at Trivandrum and Tirunelveli
26 show eastward electric field perturbations. Notably, quasi-periodic electric field fluctuations
27 with prominent periodicities of approximately 15 minutes are observed in both the IMF By and
28 ionospheric parameters, such as $dh'F/dt$ and EEJ, in both local time sectors. These fluctuations
29 are most likely driven by modulations in high-latitude disturbance polar (DP2) currents.
30 SuperDARN ionospheric convection maps display enhanced convection, rotation, and
31 expansion, which appear to influence the equatorial electric field. The presence of northward
32 Bz (NBZ) currents and reverse convection patterns likely contributes to the observed westward
33 electric field perturbations at low latitudes.

34

35 Keywords

36 Interplanetary Magnetic field (IMF) By, Prompt Penetration Electric field (PPEF),
37 Equatorial Ionosphere, Northward Bz currents (NBz), Vertical Plasma Drift.

38 1. Introduction

39 The enhanced magnetospheric convection electric field, generated by the interaction between the
40 solar wind and the magnetosphere, significantly influences ionospheric electrodynamics during
41 magnetic storms (e.g., Nishida, 1968a; 1968b; Kelley, 1979; Tsurutani et al., 2004, 2008;
42 Mannucci et al., 2005b). Under southward IMF Bz conditions, the dawn-to-dusk electric field
43 (IEFy), also known as the convection electric field, rapidly penetrates into the ionosphere. The
44 propagation of the transverse magnetic mode (TM₀) in the Earth-ionosphere waveguide plays a
45 key role in facilitating this penetration from high to low latitudes (e.g., Kikuchi et al., 1978;
46 1996). In the equatorial ionosphere, this penetrating electric field exhibits an eastward (dawn-to-
47 dusk) polarity during the day and a westward (dusk-to-dawn) polarity at night. When the IMF
48 Bz turns northward, the magnitude and direction of the convection electric field begin to weaken,
49 and an overshielding electric field develops within the plasmasphere. The polarity of this
50 overshielding field is eastward at night and westward during the day (e.g., Fejar et al., 1979;
51 Kelley et al., 1979; 2007; Kikuchi et al., 1996; Tsurutani et al., 2004, 2008; Huang, 2009, 2012;
52 Baskar and Vichare, 2013).

53 Kelley et al. [1979] introduced the concept of under- and over-shielding of the magnetospheric
54 electric field, focusing on how the Alfvén layer, a boundary region near the edge of the ring
55 current, controls the transmission of electric fields from the magnetosphere to the inner
56 magnetosphere and ionosphere. Under-shielding happens during the early stages of a
57 geomagnetic disturbance (e.g., southward turning of IMF), when the shielding currents in the
58 inner magnetosphere (especially the ring current and Region 2 field-aligned currents) have not
59 yet fully developed. In this case, magnetospheric penetration electric fields (such as those from
60 solar wind-magnetosphere interactions) can reach low-latitude ionospheric regions with little
61 attenuation, causing prompt penetration electric field (PPEF). Over-shielding, on the other
62 hand, occurs during the recovery phase of geomagnetic activity (e.g., northward turning of
63 IMF) when shielding currents still persist even after the external driver (like the solar wind
64 electric field) has decreased or stopped. In the ionosphere, this results in electric fields of
65 opposite polarity due to an overshoot of the shielding effect. Over time, studies have been built
66 on this idea, using observational data and modeling to measure the presence and effects of

67 penetration electric fields on equatorial ionospheric dynamics, including the EEJ and
68 ionospheric plasma drifts (e.g., Fejer and Scherliess 1997; Tsurutani et al., 2004, 2008; Kelley
69 et al., 2007; Huang, 2009; Singh et al., 2024, 2025; Rout et al., 2025). Huang (2009) presented
70 that the observed westward penetrating fields during substorm onset are likely due to the
71 associated overshielding effects, rather than the substorm itself. Under steady southward IMF
72 Bz conditions, substorms have been shown to induce eastward penetrating electric fields in the
73 low-latitude ionosphere (e.g., Huang, 2012). The penetration of IMF components into the
74 magnetosphere occurs under various orientations, not only during southward IMF. Tsurutani et
75 al. (1984) found that roughly 13 % of the IMF By component is transmitted into the
76 magnetotail, implying that the same physical mechanisms that drive interplanetary electric-
77 field penetration are at work.

78 Another type of penetration known as the DP2 (disturbance polar current) type penetration,
79 occurs when IMF Bz oscillates in the north and south directions. When DP2-type electric
80 field perturbations occur, the ionospheric electric field and magnetometer data reflect similar
81 types of periodic changes from high to lower latitudes (e.g., Nishida, 1968a, 1968b; Sastri
82 et al., 2000; Huang, 2019a, 2020; Singh et al., 2022; Rout et al., 2022). Aside from these
83 electric field perturbations, studies have revealed that the electric field penetration in low
84 latitude ionosphere may be triggered by several types of source mechanisms, including
85 magnetic substorms and a sudden increase in solar wind pressure (e.g., Wei et al., 2009; Rout
86 et al., 2016; Huang, 2020; Singh & Sripathi, 2021). A few studies have also shown the
87 electric field penetration associated with IMF By across equatorial and low-latitude
88 ionosphere (e.g., Kelley and Makella, 2002; Chakrabarty et al., 2017; Hui and Vichare, 2021;
89 Kumar et al., 2023). Kelley and Makella (2002) first reported the IMF By associated
90 penetration electric over low-latitude ionosphere using vertical plasma drift measurements.
91 Hui and Vichare (2021) suggested utilizing the TIEGCM model that under IMF Bz
92 southward conditions, the polarity of IMF By can considerably affect electric field
93 penetration and neutral dynamics. It is important to mention here that these studies have
94 shown the impact of the IMF By over low latitude ionosphere under the southward IMF Bz
95 conditions. Therefore, it is yet to be investigated the effects of IMF By on the low latitude
96 ionospheric electric field during the northward IMF Bz conditions.

97 This study aims to investigate the electrodynamic response of the equatorial and low-latitude
98 ionosphere to periodic changes in the IMF By component under conditions of steady northward
99 IMF Bz. It is important to highlight that previous research has primarily focused on the effects

100 of IMF B_y during southward IMF B_z conditions, which are typically associated with enhanced
101 magnetospheric convection and prompt penetration electric fields (e.g., Kelley and Makela,
102 2002; Chakrabarty et al., 2017). However, to the best of the authors' knowledge, no systematic
103 study has been conducted to investigate how periodic oscillations in IMF B_y influence the
104 penetration of electric fields over the low-latitude ionosphere during steady northward IMF B_z
105 conditions. This study provides the first observational evidence of ionospheric disturbances
106 linked to penetrating electric fields driven by IMF B_y fluctuations during sustained northward
107 IMF B_z . Through the combination of measurements from opposite longitude sectors (both from
108 India and American longitude), the study aims to reveal new insights into the magnetosphere-
109 ionosphere coupling mechanism.

110 The structure of this study is as follows: Section 2 describes the source of data analysis.
111 Observations and results are presented in Section 3. Section 3.1 contains the solar
112 wind geomagnetic conditions, and Section 3.2 presents the ground-based measurements.
113 Sections 4 and 5 provide the discussion and conclusion.

114 **2. Data Sets**

115 In this study, we used digisonde and Canadian Advanced Digital Ionosonde (CADI) data
116 to calculate the F-layer height ($h'F$) and rate change of height ($dh'F/dt$) at Jicamarca (JIC:
117 11.9°S; -76.8°E), Fortaleza (FTZ: 3.9°S, -38.52°E), Trivandrum (TVM: 8.5°N; 77°E),
118 Tirunelveli (TIR: 8.73° N, 77.70° E), and Guam (GUA: 13.69°N; 144.87°E) stations. The
119 Jicamarca Incoherent Scatter Radar (ISR) recorded vertical plasma drift measured by the
120 JULIA MP (Jicamarca Unattended Long-Term Investigations of the Ionosphere and
121 Atmosphere Medium Power) mode was also used. The time interval for JULIA radar data
122 was 5 minutes, whereas the ionograms on the FTZ and JIC were recorded at 10- and
123 15-minute intervals, respectively. The ionogram intervals at TVM, TIR, and GUA were 7,
124 10, and 15 minutes, respectively. The local times at the locations are given by $LT = UT + 5.5$
125 hours (at TVM and TIR), $LT = UT + 9.6$ hours (at GUA), $LT = UT - 5$ hours (at JIC), and LT
126 $= UT - 3$ hours (at FTZ). The two magnetometers located at equatorial and off-equatorial sites
127 at Jicamarca (1.9°S, 76.8°W, $\sim 0.5^\circ$ N dip latitude) and Piura (5.2°S, 80.6°W, $\sim 6.81^\circ$ N dip
128 latitude) are used to calculate the EEJ. The EEJ is a key feature of low-latitude electrodynamics.
129 It is influenced by both external factors, such as IMF-driven electric fields, and internal drivers
130 like tidal winds and conductivity variations (e.g., Zhang et al., 2021). The EEJ consists of a
131 narrow band of intense eastward current flowing in the E-region ionosphere. It is typically
132 confined within $\pm 3^\circ$ magnetic latitude of the magnetic equator and occurs at an altitude of

133 approximately 100–110 km. The solar wind parameters are obtained from ACE satellite
134 measurements, and the geomagnetic indices are sourced from WDC Kyoto. A global network
135 of real-time magnetic observatory magnetometers is used to acquire magnetic field data.
136 SuperDARN observations are used to generate convection maps of the Northern Hemisphere
137 at high latitudes.

138 **3. Observations and Results**

139 **3.1 Solar Wind and Geomagnetic Conditions**

140 The interplanetary and geomagnetic conditions on December 19, 2015, between 15:00 and
141 20:00 UT are presented in Figure 1. The change in solar wind speed (V_{sw} , red), solar wind
142 dynamic pressure (P_{dyn} , black), and particle density (N_p , green) are shown in Figure 1a.
143 Figures 1b-1h show variations in IMF B_x , B_y , and B_z , IEF $_y$, auroral electrojets (AU and
144 AL), and Sym-H. Except for the IMF B_y component interplanetary and geomagnetic
145 conditions, no significant changes were exhibited between 16:30 and 18:00 UT, as shown in
146 the grey shaded area. Interestingly, under steady northward B_z ($\sim +15$ nT) conditions, and with
147 stable values of V_{sw} , N_p , and P_{dyn} at 450 km/s, 20 cm^{-3} , and 8 nPa, respectively, the IMF B_y
148 fluctuated between ± 8 nT, showing three significant peaks. Positive and negative peaks were
149 noticed at 16:55 and 17:06 UT, 17:12 and 17:20 UT, and 17:27 and 17:35 UT, respectively. At
150 the same time, Sym-H expressed an eastward ring current with a value of nearly +10 nT. The
151 positive and negative values of B_y correspond to the dusk and dawn side polarities of the
152 magnetic field. It must be noted that this observation is unique as it has been found after
153 carefully analyzing more than 30 years of solar wind data. Figure 2 shows the changes in
154 energetic particle fluxes (e.g., electron flux) between 15:00 and 20:00 UT. The electron fluxes
155 were detected at different energy levels (i.e., 40, 75, 150, 275, 475 keV) by nine telescopes
156 (numbered 101-109) on geosynchronous satellites (GOES-13 and 15). The left panels show
157 observations from GOES-13, while the right panels display observations from GOES-15. The
158 magnetic footpoints of GOES-13 and 15 were located at Geog: 57.5°N , 79.2°W and 59.5°N ,
159 129.2°W , respectively. Dispersionless particle injection at geosynchronous was not observed
160 during the noon (between 01:34 and 12:56 MLT) and morning (between 07:35 and 08:55
161 MLT) hours. The GOES 13 shows a dispersion less particle injection at 18:30-19:00 UT
162 (09:25-09:55 MLT), but the location and local time duration do not coincide with our results.
163 However, AU/AL (Figure 1e) also did not observe any substorm signatures; the initial increase
164 in particle injection at 16:15 UT (Figure 2) appears to be caused by a rise in solar dynamic
165 pressure (Figure 1e).

166 3.2 Ground-based Ionosonde and Radar Observations over Indian and American 167 Longitudes

168 Figure 3 shows ground-based ionosonde observations over the Indian longitudes during
169 nighttime hours between 15:30 and 19:00 UT (21:00 and 00:30 LT). Figures 3a, 3b, and 3c
170 illustrate the virtual height ($h'F$) variations, as well as the mean and standard deviation
171 estimated from five International Quiet Days (IQDs) for Guam, Tirunelveli, and Trivandrum,
172 respectively. Figure 3d depicts the vertical plasma drift ($dh'F/dt$), which is derived using the
173 rate of change in ionospheric height (i.e., F-layer height) at Trivandrum. Figure 3e lists
174 variations in IMF B_y (blue) and B_z (red). Interestingly, the $h'F$ and $dh'F/dt$ variations shown
175 in the vertical cyan shaded regions coincide well with IMF B_y changes and exhibit quasi-
176 periodic oscillations, similar to IMF B_y between 16:45-18:00 UT. The time resolution and data
177 intervals are not the same in each observation. Ionosonde observations from Tirunelveli and
178 Guam stations are used to validate Trivandrum observations. The $h'F$ enhancements and
179 oscillations in virtual height and vertical plasma drift within the shaded region could be
180 attributed to disturbances in the eastward electric field. Notably, the vertical plasma drift and
181 virtual height at the dip equator indicate the presence of an eastward electric field rather than
182 the typically expected westward field during nighttime hours.

183 Figure 4 presents ionosondes, ISR radar, and magnetometers observations during daytime
184 hours over the American longitude sector, along with IMF data. Figure 4a and 4b show
185 variations in the ionospheric virtual height ($h'F$) along with the five-day mean and standard
186 deviation of IQD variations at the Jicamarca and Fortaleza stations, respectively. Figure 4c
187 presents the vertical plasma drift (V_d) measured by the Jicamarca ISR operating in JULIA MP
188 mode (black circles), along with the drift predicted by the Scherliess and Fejer (SF) model,
189 represented by the black line (e.g., Fejer & Scherliess, 1997). Figure 4d presents the EEJ
190 together with its five-day IQD mean and standard deviation. Figure 4e displays variations in
191 the IMF components B_y (blue) and B_z (red), as well as the detrended ΔEEJ (green). The ΔEEJ
192 is calculated by subtracting the mean EEJ values from the event-specific EEJ values. The
193 vertical plasma drift, EEJ, and virtual height exhibited a sudden increase caused by the
194 eastward enhanced electric field generated by the southward IMF B_z at $\sim 16:20$ UT. Afterward,
195 the ISR drift and EEJ showed a steady reduction in their values due to the westward electric
196 field between 16:35 and 18:05 UT, corresponding to a steady northward IMF B_z . However,
197 during this constant reduction, the ISR drift and EEJ exhibited oscillations corresponding to
198 the IMF B_y . Subsequently, a sudden increase was observed in the ISR drift, EEJ, and $h'F$,

199 corresponding to the southward IMF Bz between 18:00 and 18:30 UT.
200 Notably, the vertical plasma drift, EEJ, and h'F (at Jicamarca) exhibit a significant reduction
201 compared to the quiet-time mean variations. This reduction is attributed to the westward prompt
202 penetration electric field, which opposes the background eastward field, leading to the observed
203 decrease between 16:35 and 18:00 UT (Figures 4b, 4c, and 4d). On the dayside, the eastward
204 penetration electric fields are not strong enough to fully cancel the westward electric field,
205 allowing the westward electric field oscillations to persist. The ionosonde measurements from
206 Fortaleza further confirm the electric field disturbance, as significant oscillations in h'F are
207 evident (Figure 4a). The clear oscillations in the westward electric field disturbance are shown
208 in Δ EEJ, which closely follow the IMF By oscillations (Figure 4e). These findings support a
209 correlation between IMF By variations and the oscillations observed in the EEJ, vertical plasma
210 drift, and h'F, driven by IMF By. The disturbed electric field oscillations associated with IMF
211 By fluctuations are highlighted in the cyan-shaded region.

212 To understand the ionospheric perturbations in the electric field carried by IMF By
213 reorientations, Figure 5 illustrates a periodogram analysis using the Lomb-Scargle (LS)
214 method and Morlet wavelet analysis (e.g., Lomb, 1976; Torrence and Compo, 1998) for IMF
215 By, EEJ, and dh'F/dt at Trivandrum. The left panels show the LS periodogram, the middle
216 panels show the wavelet spectrogram, and the right panels show the global wavelet spectrum
217 (GWS). The third-order Savitzky-Golay smoothing algorithm was used to eliminate the
218 long-term fluctuations for the periodogram analysis (e.g., Savitzky & Golay, 1964; Singh et
219 al., 2022). The LS periodogram analysis and GWS reveal a dominant common periodicity
220 of ~ 15 minutes in the ionospheric parameters, as well as in the IMF By. Wavelet analysis
221 demonstrated spectral power enhancements centred at periods of around 15 minutes and
222 showed common periodicity. According to this analysis, the reorientations of the IMF By
223 by carrying out electric field circumstances are substantially triggering the common
224 periodic oscillations in the ionosphere.

225 Figure 6 provides the DP2 contours and convection maps at the following times: (a) 17:00-
226 17:02 and 17:06-17:08 UT; (b) 17:18-17:20 and 17:20-17:22 UT; and (c) 17:26- 17:28 and
227 17:30-17:32 UT. The contour maps with polar cap potentials are shown in the supporting
228 Figure (S). From the Figures, it can be seen that reverse convection or NBZ (northward
229 Bz-component) currents exist when the IMF By polarity is negative or in the dusk direction
230 (Figure 6a, 6c, and S); during the IMF By positive condition, the NBZ currents are
231 very weak. The convection cell grew larger toward lower latitudes and rotated toward

232 the dawnside during the time of positive IMF B_y (Figure 7b and S).

233 **4. Discussion**

234 The region-1 (R1) and region-2 (R2) field-aligned currents (FACs) system and related
235 magnetospheric plasma convection are driven by the dynamic interaction of the solar wind
236 and the magnetosphere. R1, R2, and their horizontal currents produce the DP2 (disturbance
237 polar current) current system at high latitudes, create an extra electric field there, and undergo
238 large changes with changes in the IMF B_z (e.g., Kikuchi et al., 1978; 1996). In the plasma
239 sheet, the magnetospheric convection electric field (dawn to dusk potential drop) gets stronger
240 under the southward IMF B_z condition. This enhanced convection electric field
241 subsequently interacts with the high-latitude ionosphere, causes disruption to the DP2-type
242 current system, and promotes the DP2-type current system to reach lower latitudes instantly
243 (e.g., Nishida, 1968a, 1968b; Kikuchi et al., 1996). In the local noon hours, an additional
244 current system that exists just poleward of the R1 FACs occurs at times of almost entirely
245 northward IMF B_z and is referred to as NBZ (northward B_z) or cusp currents (e.g.,
246 Iijima and Potemra, 1978), or sunward convection (e.g., Vasyliunas, 1989; Vennerstrom et
247 al., 2002). The NBZ current system is composed of two localized FAC regions that are
248 symmetrically positioned at local noon in both hemispheres and move in the opposite direction
249 compared to R1 FACs, which enter and exit the ionosphere on the dawn and dusk sectors. The
250 IMF B_y causes large asymmetries in the magnetosphere that are reflected in global currents,
251 convection patterns, and auroral intensity by changing the flow of plasma in both hemispheres
252 in various ways (e.g., Weimer, 1995, 1996). In our observations on the dayside, the eastward
253 enhanced electric field was caused by the southward turning of IMF B_z (e.g., Tsurutani et al.,
254 2004; 2008; Fejar et al., 1979), which enhanced the EEJ, increased the vertical plasma drift,
255 and raised the F-layer height over the American longitudes, as shown in Figure 4 at around
256 16:20 and 18:05 UT (e.g., Singh et al., 2024, 2025; Rout et al., 2025). At the same time, on the
257 nightside over Indian longitudes, a reduction in $h'F$ was observed, which was caused by the
258 westward penetration electric field (Figure 3).

259 Figures 3 and 4 exhibited eastward (on the nightside, between 21:30 and 23:00 LT) and
260 westward (on the dayside, between 11:30 and 13:00 LT) electric field penetration signatures
261 over Indian and American longitudes, respectively. These penetration electric fields with
262 opposite polarity are caused by reverse convection or NBZ currents (e.g., Huang, 2019a).
263 In our observations, Figures 6a and 6c demonstrate the formation of reverse convection and
264 NBZ currents. Huang (2019a) observed a westward penetration electric field on the dayside

265 and suggested that reverse convection and NBZ currents are the main generators for this
266 electric field. Figure 6 displays the ionospheric convection maps in the northern hemisphere
267 made by the SuperDARN of HF (high frequency) radars between 17:00 and 17:32 UT
268 under the constant northward IMF B_z condition, attempting to understand the electric field
269 penetration driven by IMF B_y . Arrows in the SuperDARN convection maps indicate
270 convection vectors computed from the radar echoes, while contours in the maps display
271 statistical data. The convection maps show the occurrence of reverse convection (sunward)
272 or NBZ currents within the polar region. Figures 6a and 6c show ionospheric convection maps
273 at 17:00, 17:06, 17:26, and 17:32 UT when IMF B_y turns from positive to negative (dawn to
274 dusk) direction or is in duskside (e.g., negative B_y). Most importantly, as the reverse
275 convection occurs, dusk and dawn cells split into multiple cells and shift towards lower
276 latitudes. The westward and eastward electric field disturbances are observed at the equator
277 on the dayside and nightside. These penetrating electric fields are generated by reverse
278 convection or NBZ current as reported by Huang (2019a).

279 According to Heelis (1984), significant changes in the interplanetary magnetic field (IMF) B_y
280 can effectively modify the DP2 currents, causing them to expand toward lower latitudes. The
281 size and morphology of the convection cells in the morning and evening sectors are controlled
282 by both the magnitude and direction of IMF B_y and B_z , as well as by distortions in the DP2
283 current pattern induced by IMF B_y fluctuations (e.g., Heelis et al., 1982; Heelis, 1984;
284 Ruohoniemi and Greenwald, 2005; Tenfjord et al., 2015). Previous studies have investigated
285 how the IMF influences high-latitude plasma convection and DP2 currents (e.g., Heelis, 1984;
286 Ruohoniemi and Greenwald, 2005; Tanaka, 2007).

287 In our observations, under positive IMF B_y (i.e., duskward B_y perturbations), the dusk cell
288 extended toward the dayside and adopted an approximately circular shape, whereas the dawn
289 cell assumed a crescent shape (Figure 6b). Simultaneously, the overall convection pattern
290 rotated clockwise about the noon midnight meridian, yielding anticlockwise plasma flow in the
291 nightside sector and a duskward zonal flow near midnight. Conversely, for negative IMF B_y
292 (i.e., dawnward B_y perturbations), the dawn cell became nearly circular and the dusk cell
293 became crescent-shaped; the entire pattern again rotated clockwise, producing dawnward zonal
294 flow at midnight (Figures 6a, 6c). These distortions in the polar convection cells appear to be
295 regulated by fluctuations in IMF B_y under sustained northward B_z conditions, thereby
296 triggering oscillations in the equatorial electric fields.

297 The periodic oscillations of disturbance electric fields, driven by various solar wind

298 magnetosphere interaction processes, have been extensively investigated in earlier studies (e.g.,
299 Nishida, 1968a, 1968b; Sastri et al., 2000; Huang, 2019a; Singh et al., 2022; Rout et al., 2022).
300 Near the magnetic equator, Nishida (1968a, 1968b) reported quasi-periodic oscillations in the
301 geomagnetic field recorded by ground-based magnetometers, which coincided with IMF Bz
302 fluctuations. Subsequent studies have identified a broad range of periodicities associated with
303 DP2-type electric field fluctuations, including 25–35 minutes (e.g., Sastri et al., 2000), 30–40
304 minutes (e.g., Kikuchi et al., 1996; Chakrabarty et al., 2015), 40 minutes (e.g., Raut et al., 2019),
305 and dual periodicities of 30 and 60 minutes (e.g., Rout et al., 2017; Singh et al., 2022). In this
306 study, we report for the first time a DP2-type fluctuation with a 15-minute periodicity, which
307 appears to be linked to fluctuations in the IMF By component. Figure 5 reveals a consistent and
308 dominant 15-minute periodicity in IMF By, EEJ, and dF/dt at both longitudes. These shared
309 periodicities between IMF By and ionospheric parameters suggest that DP2-associated
310 perturbations influence the equatorial electric field. Owing to the rotation and expansion of the
311 dawnside convection cell, both the dayside (Jicamarca) and nightside (Trivandrum) stations
312 observed these electric field oscillations, indicating a longitudinally coherent response (e.g.,
313 Chakrabarty et al., 2017).

314 5. Summary

315 This study highlights the dominant role of the IMF By in causing the electric field disturbances
316 in the ionospheric system at opposite longitude sectors at the same time. During northward IMF
317 Bz, the westward electric field in the dayside equatorial zone should be seen as coming
318 from two separate sources: reverse convection (e.g., Huang et al., 2019a) in the polar region
319 and R-2 FACs (e.g., Kikuchi et al., 1978; 1996). Conversely, the polar reverse convection
320 can be responsible for transmitting the westward electric field. In our case, the opposite
321 polarity of the disturbed electric fields observed at equatorial latitudes is most likely driven by
322 NBZ currents, rather than by R2 FACs, which typically occur at much lower latitudes. However,
323 the similar oscillatory behavior in the disturbed electric fields appears to result from the combined
324 effects of strong convection, enlargement, and rotations of the convection cells. Under constant
325 northward IMF Bz conditions, a DP2-type prompt penetration electric field of 15-minute
326 periodicity is found to be observed in IMF By, ionospheric vertical plasma drift, and the EEJ.
327 The southward turning of the IMF Bz heightened the dayside eastward electric field, which
328 strengthened the EEJ, increased vertical plasma drift, F-region virtual height, and the nightside
329 westward electric field disturbance.

330

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333 Research Organization (ISRO) and finalized at the Jicamarca Radio Observatory in Lima, Peru.
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344 **Declaration of Competing Interest**

345 The authors declare that they have no known competing financial interests or personal
346 relationships that could have appeared to influence the work reported in this paper.

347 **Data Availability Statement**

348 The Jicamarca Radio Observatory's Madrigal Database provided the JULIA radar data used in
349 this study (<http://jro1.igp.gob.pe/madrigal/>). The Jicamarca Radio Observatory is a facility run
350 by the Instituto Geofísico del Peru with funding from the NSFAGS-1433968 through Cornell
351 University. Data for the interplanetary and geomagnetic conditions were obtained from the
352 NASA GSFC CDAWeb (http://cdaweb.gsfc.nasa.gov/istp_public/) and WDC kyoto
353 (<http://wdc.kugi.kyoto-u.ac.jp>). The international real-time magnetic observatory
354 magnetometer network (https://imag-data.bgs.ac.uk/GIN_V1/GINForms2) was used to gather
355 magnetic field information. The convection maps were downloaded from
356 SuperDARN Network (<https://superdarn.ca/>). The Global Ionosphere Radio Observatory
357 (GIRO) online provides access to the ionosonde data at <https://giro.uml.edu/didbase/>.

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494 **Figures**

495 **Figure 1.** Solar wind and geomagnetic conditions on 19 December 2015. (a) X-component of
 496 the interplanetary magnetic field (IMF Bx); (b) Y- and Z-components of the IMF (By and Bz,

497 in nT); (c) Y-component of the interplanetary electric field (IEF E_y , in mV/m); (d) solar wind
 498 dynamic pressure (P_{dyn} , in nPa), proton density (N_p , in cm^{-3}), and solar wind speed (V_{sw} , in
 499 km/s); (e) auroral electrojet indices (AU and AL, in nT); and (f) Sym-H index (in nT). The
 500 grey-shaded region highlights significant changes in the IMF Y-component between 16:30 and
 501 18:00 UT.

502 **Figure 2.** Electron flux observations from the geosynchronous GOES satellites across multiple
 503 telescopes and energy channels. (a) GOES-13 and (b) GOES-15. The grey-shaded area
 504 indicates the time interval during which no significant changes in electron flux were observed.

505 **Figure 3.** Variations in the virtual height of the F-region ($h'F$) at (a) Guam (GUA), (b)
 506 Tirunelveli (TIR), and (c) Trivandrum (TVM); (d) vertical plasma drift ($dh'F/dt$) derived from
 507 $h'F$ variations; (e) IMF B_z and B_y components. The grey lines with error bars represent the
 508 mean and standard deviation from five International Quiet Days (IQDs). The cyan-shaded
 509 region marks the period of significant ionospheric perturbations

510 **Figure 4.** (a–b) Variations in $h'F$ at Fortaleza and Jicamarca; (c) vertical plasma drift from the
 511 Jicamarca ISR (JULIA, MP mode, shown as black circles) along with the Scherliess-Fejer (SF)
 512 model drift (black line); (d) equatorial electrojet (EEJ, green) along with the IQD mean and
 513 standard deviation (grey); and (e) IMF B_z (red), B_y (blue), and the detrended ΔEEJ (green).
 514 The grey-shaded region and vertical lines indicate periods of significant disturbances.

515 **Figure 5.** Periodogram analysis of IMF B_y , EEJ, and vertical plasma drift ($dh'F/dt$) over the
 516 Indian and American sectors. Left panels show the Lomb-Scargle (L-S) periodogram analysis,
 517 middle panels display the Morlet wavelet spectra, and right panels present the corresponding
 518 global wavelet spectra (GW).

519 **Figure 6.** SuperDARN ionospheric convection maps along with DP2 contours over the
 520 northern hemisphere on 19 December 2015, at (a) 17:00-17:02 and 17:06-17:08 UT (reverse
 521 convection; IMF B_y negative); (b) 17:18-17:20 and 17:20-17:22 UT (dusk cell
 522 rotation/expansion; IMF B_y positive); and (c) 17:26-17:28 and 17:30-17:32 UT (reverse
 523 convection; IMF B_y negative).

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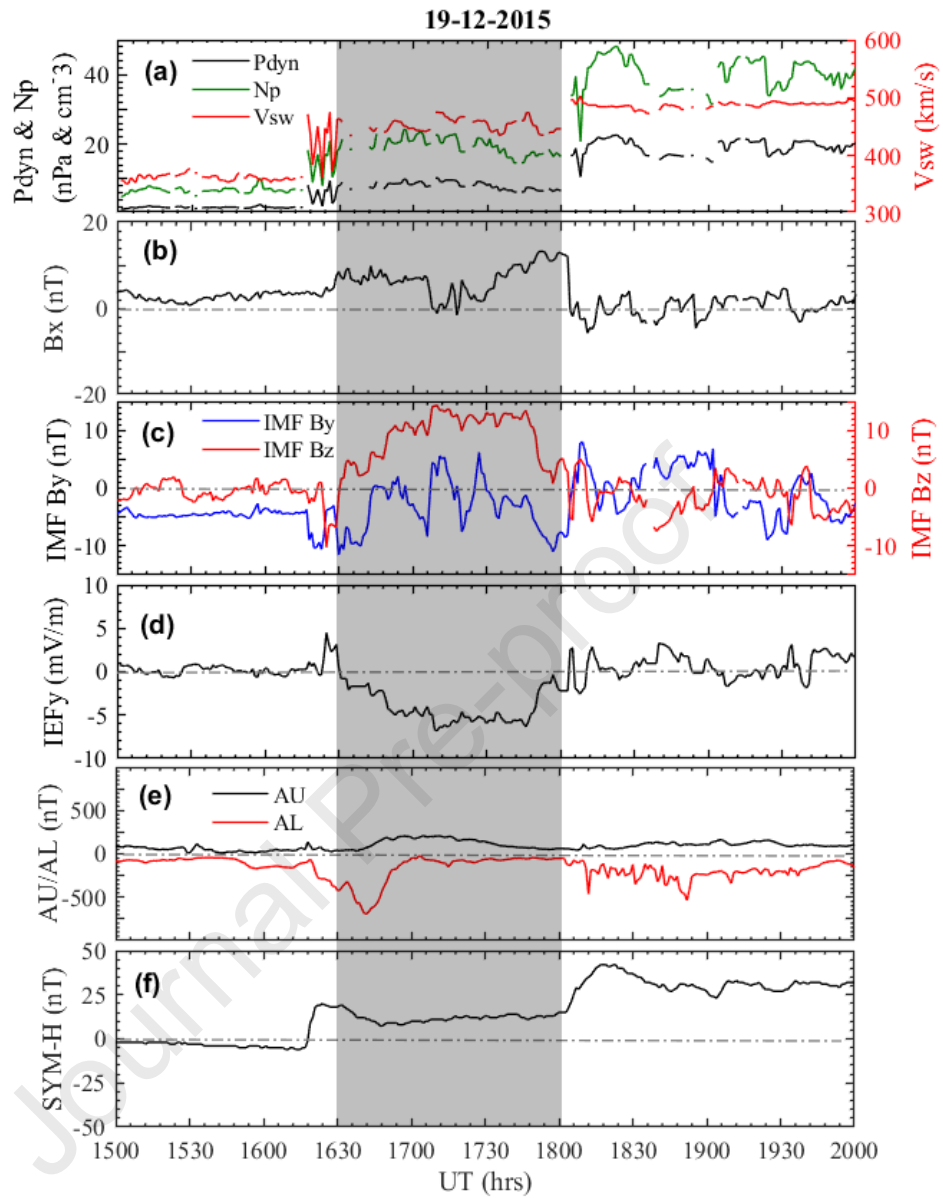
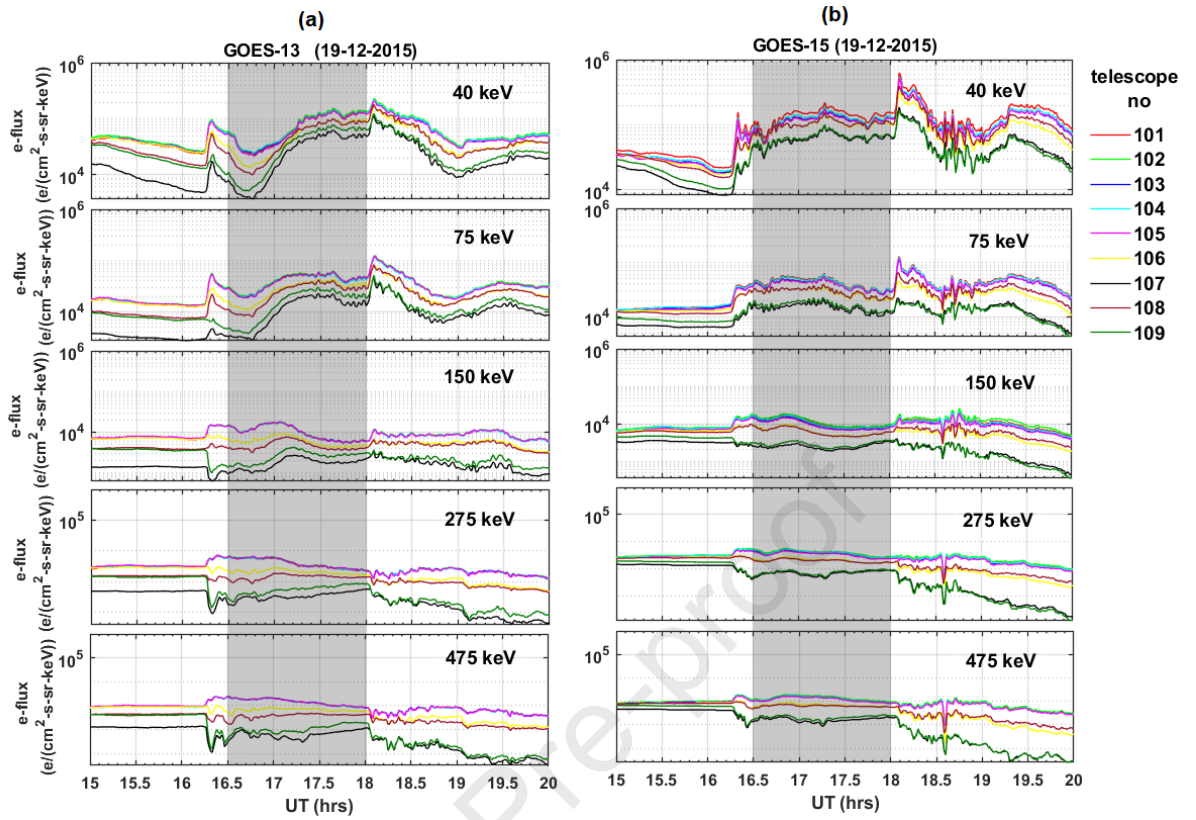


Figure 1

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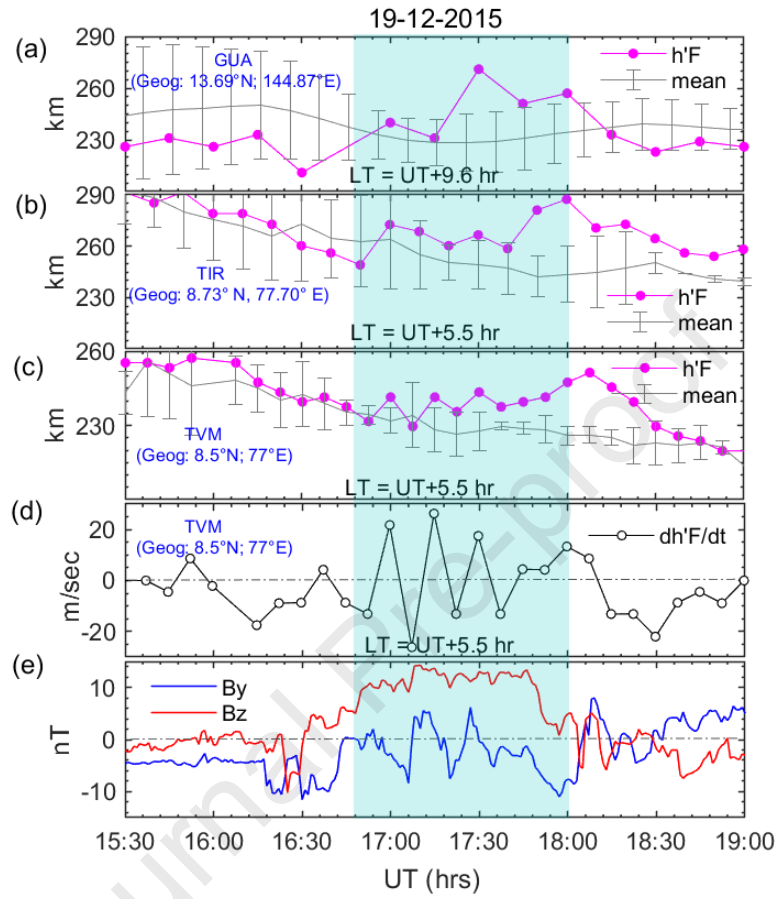
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Figure 2

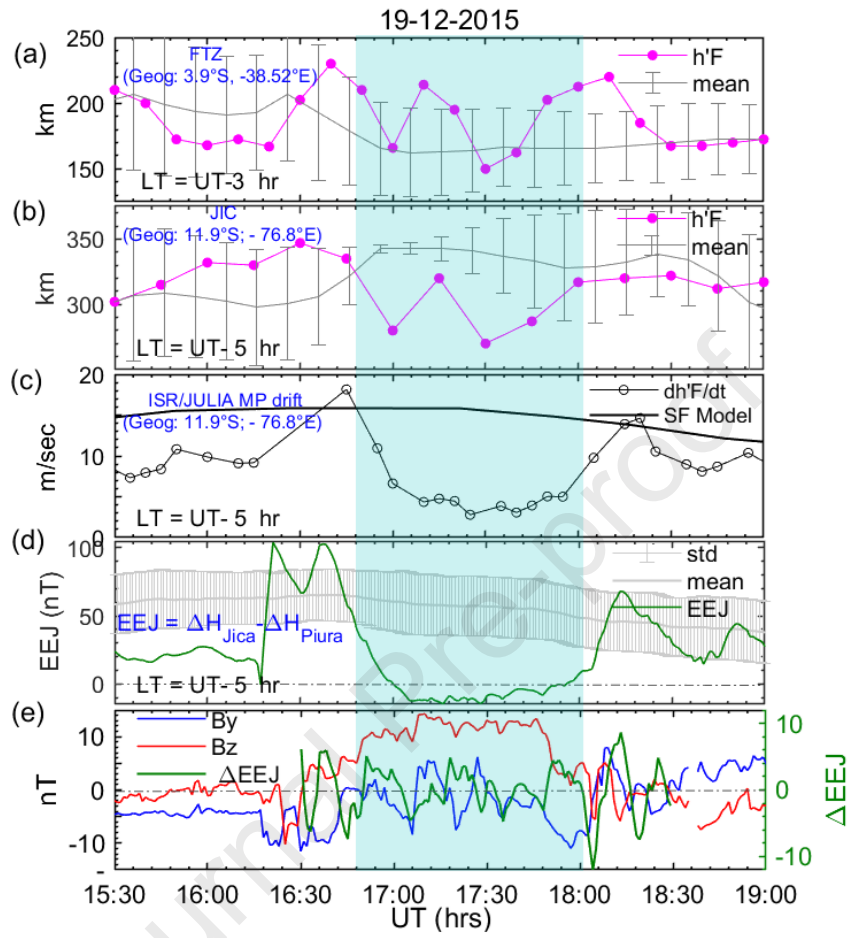
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Figure 3

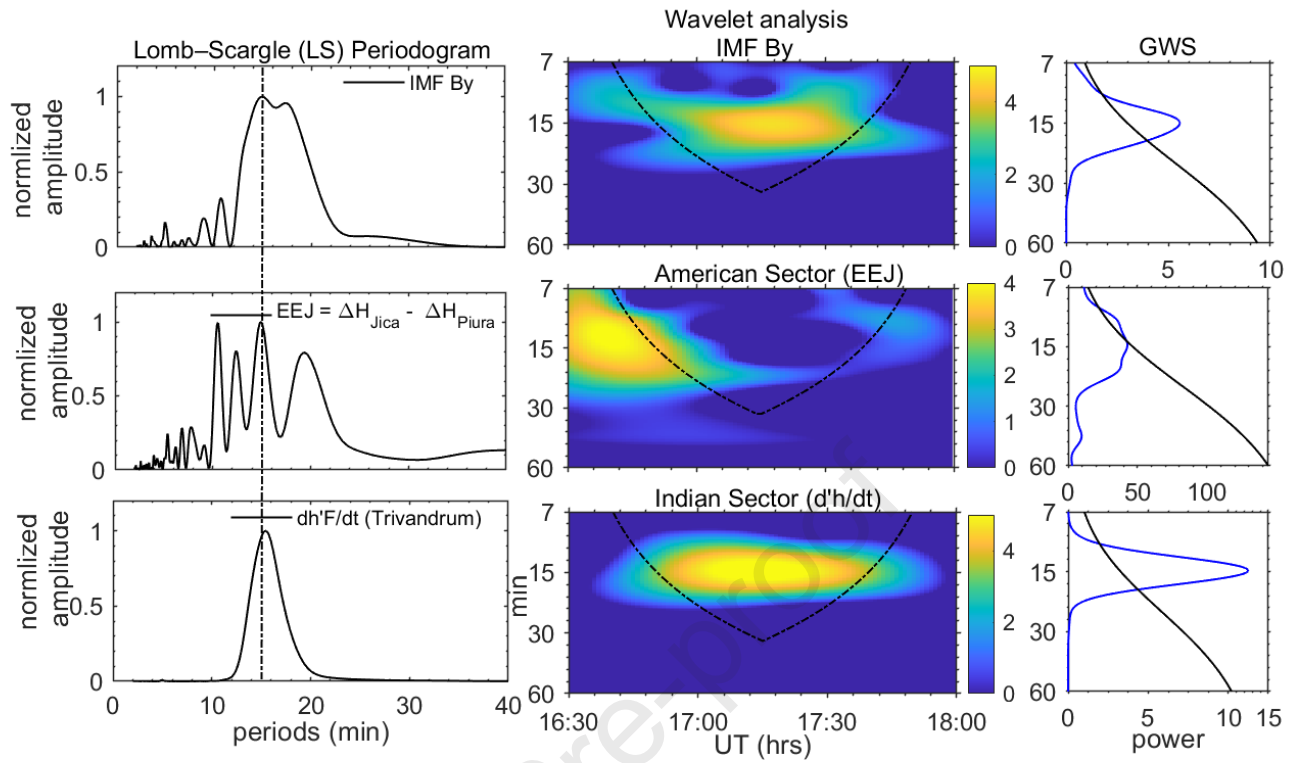
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Figure 4

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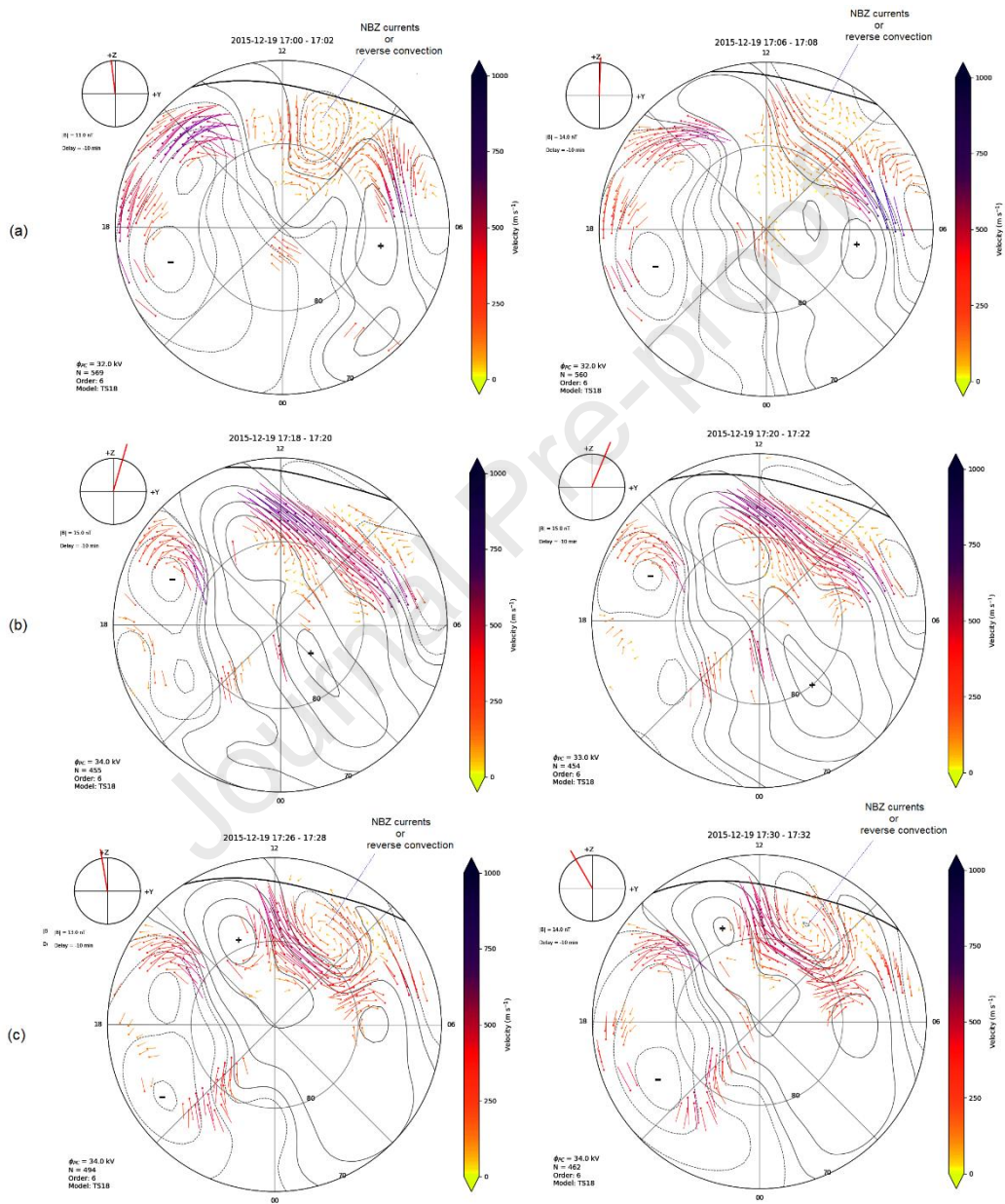
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Figure 5

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Figure 6

1. Contrary to expectations, the same polarity penetration electric field observed at opposite longitudes during day and night local times.
2. IMF B_y -related quasi-periodic (~15 minutes) signature of penetration electric field.
3. Reverse convection and NBZ (northward B_z) currents cause an electric field with opposite polarity penetration to exist.

Declaration of Interest Statement

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper